THE BERWALD-MOOR METRIC IN THE TANGENT BUNDLE OF THE SECOND ORDER

Gheorghe Atanasiu and Nicoleta Brinzei

Department of Algebra and Geometry, Transilvania University, Brasov, Romania gh atanasiu@yahoo.com, nico.brinzei@rdslink.ro

As an application of the results of the first author obtained in the papers [1] and [2], the geometry of the second order tangent bundle T^2M (or second order jet bundle J_0^2M) endowed with two special types of metrics compatible with the 2-contact structures is studied. The particularity of these two models is that the horizontal and the $v^{(1)}$ - part of the metric are both given by the same Riemannian metric (respectively, its horizontal part is Riemannian), while its $v^{(2)}$ -part is given by the flag-Finsler Berwald-Moor metric (respectively, the $v^{(1)}$ and $v^{(2)}$ - parts are given by the flag-Finsler Berwald-Moor metric, [5]).

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1 The 2-Tangent Bundle T^2M

Let M be a real 4-dimensional manifold of class \mathcal{C}^{∞} , (T^2M, π^2, M) its second order tangent bundle, [1], and let $\widetilde{T^2M}$ be the space T^2M without its null section. For a point $u \in T^2M$, let $(x^i, y^{(1)i}, y^{(2)i})$ be its coordinates in a local chart.

Let N be a nonlinear connection, [3], [8]- [13], and denote its coefficients by $\binom{N_i^i, N_j^i}{1^j, 2^j}$, i, j = 1, ..., 4. Then, N determines the direct decomposition

$$T_u T^2 M = N_0(u) \oplus N_1(u) \oplus V_2(u), \ \forall u \in T^2 M.$$
(1)

The adapted basis to (1) is $(\delta_i, \delta_{1i}, \delta_{2i})$ and its dual basis is $(dx^i, \delta y^{(1)i}, \delta y^{(2)i})$, where

$$\begin{cases}
\delta_{i} = \frac{\delta}{\delta x^{i}} = \frac{\partial}{\partial x^{i}} - N_{1}^{k} \frac{\partial}{\partial y^{(1)k}} - N_{2}^{k} \frac{\partial}{\partial y^{(2)k}} \\
\delta_{1i} = \frac{\delta}{\delta y^{(1)i}} = \frac{\partial}{\partial y^{(1)i}} - N_{1}^{k} \frac{\partial}{\partial y^{(2)k}} \\
\delta_{2i} = \frac{\partial}{\partial y^{(2)i}} = \dot{\partial}_{2i},
\end{cases}$$
(2)

respectively,

$$\begin{cases} \delta y^{(1)i} = dy^{(1)i} + M_k^i dx^k \\ \delta y^{(2)i} = dy^{(2)i} + M_k^i dy^{(1)k} + M_k^i dx^k, \end{cases}$$
(3)

where M_{1k}^i, M_{k}^i are the dual coefficients of the nonlinear connection N.

Then, a vector field $X \in \mathcal{X}(T^2M)$ is represented in the local adapted basis as

$$X = X^{(0)i}\delta_i + X^{(1)i}\delta_{1i} + X^{(2)i}\delta_{2i},$$
(4)

with the three right terms (called *d*-vector fields) belonging to the distributions N, N_1 and V_2 respectively.

A 1-form $\omega \in \mathcal{X}^*(T^2M)$ will be decomposed as

$$\omega = \omega_i^{(0)} dx^i + \omega_i^{(1)} \delta y^{(1)i} + \omega_i^{(2)} \delta y^{(2)i}.$$

Similarly, a tensor field $T \in \mathcal{T}_s^r(T^2M)$ can be split with respect to (1) into components, which will be called *d*-tensor fields.

2 N-linear connections. d-tensors of curvature

An *N*-linear connection D, [1], [2], is a linear connection on T^2M , which preserves by parallelism the distributions N, N_1 and V_2 .

An N-linear connection is locally given by its coefficients

$$D\Gamma(N) = \left(\begin{array}{c} L^{i}_{jk}, L^{i}_{(10)}, L^{i}_{jk}, L^{i}_{(20)}, L^{i}_{jk}, C^{i}_{(11)}, L^{i}_{jk}, C^{i}_{(21)}, L^{i}_{jk}, C^{i}_{(22)}, L^{i}_{jk}, C^{i}_{(22)}, L^{i}_{jk}, C^{i}_{(22)}, L^{i}_{jk}, L^{i}_{(22)}, L^{i}_{(22)}, L^{i}_{jk}, L^{i}_{(22)}, L^{i}_{(22)}$$

where

$$\begin{cases}
D_{\delta_k}\delta_j = \sum_{(00)}^{i}{}_{jk}i_i, D_{\delta_k}\delta_{1j} = \sum_{(10)}^{i}{}_{jk}i_i, D_{\delta_k}\delta_{2j} = \sum_{(20)}^{i}{}_{jk}i_k\delta_{2i} \\
D_{\delta_{1k}}\delta_j = \sum_{(01)}^{i}{}_{jk}i_i, D_{\delta_{1k}}\delta_{1j} = \sum_{(11)}^{i}{}_{jk}i_k\delta_{1i}, D_{\delta_{1k}}\delta_{2j} = \sum_{(21)}^{i}{}_{jk}i_k\delta_{2i} \\
D_{\delta_{2k}}\delta_j = \sum_{(02)}^{i}{}_{jk}i_k\delta_i, D_{\delta_{2k}}\delta_{1j} = \sum_{(12)}^{i}{}_{jk}i_k\delta_{1i}, D_{\delta_{2k}}\delta_{2j} = \sum_{(22)}^{i}{}_{jk}i_k\delta_{2i}
\end{cases}$$
(6)

The curvature of the N-linear connection D,

$$R(X,Y)Z = D_X D_Y Z - D_Y D_X Z - D_{[X,Y]}Z,$$

is completely determined by its components (which are *d*-tensors) $R(\delta_{\gamma l}, \delta_{\beta k}) \delta_{\alpha j}$. Namely, the 2-forms of curvature of an *N*- linear connection are, [1], [2],

$$\Omega_{(\alpha)}^{i}{}_{j} = \frac{1}{2} \frac{R}{(0\alpha)}{}_{jkl}^{i} dx^{k} \wedge dx^{l} + \frac{P}{(1\alpha)}{}_{jkl}^{i} dx^{k} \wedge \delta y^{(1)l} + \\
+ \frac{P}{(2\alpha)}{}_{jkl}^{i} dx^{k} \wedge \delta y^{(2)l} + \frac{1}{2} \frac{S}{(1\alpha)}{}_{jkl}^{i} \delta y^{(1)k} \wedge \delta y^{(1)l} + \\
+ \frac{Q}{(2\alpha)}{}_{jkl}^{i} \delta y^{(1)k} \wedge \delta y^{(2)l} + \frac{1}{2} \frac{S}{(2\alpha)}{}_{jkl}^{i} \delta y^{(2)k} \wedge \delta y^{(2)l},$$
(7)

 $\alpha = 0, 1, 2$, where the coefficients $\underset{(0\alpha)}{R}_{jkl}^{i}, \underset{(\beta\alpha)}{P}_{jkl}^{i}, \underset{(2\alpha)}{Q}_{jkl}^{i}, \underset{(\beta\alpha)}{S}_{jkl}^{i}$ are d-tensors, named the *d*-tensors of curvature of the N-linear connection D.

3 Metric structures on T^2M

A Riemannian metric on T^2M is a tensor field G of type (0, 2), which is nondegenerate in each $u \in T^2M$ and positively defined on T^2M .

In this paper, we shall consider only metrics in the form

$$G = g_{ij} dx^{i} \otimes dx^{j} + g_{ij} \delta y^{(1)i} \otimes \delta y^{(1)j} + g_{ij} \delta y^{(2)i} \otimes \delta y^{(2)j},$$
(8)

where $g_{ij} = g_{ij}(x, y^{(1)}, y^{(2)})$; this is, such that the distributions N, N_1 and V_2 generated by the nonlinear connection N be orthogonal in pairs with respect to G.

Let also

$$F = \sqrt[4]{y^{(1)1}y^{(1)2}y^{(1)3}y^{(1)4}}$$

be the Berwald-Moor Finsler function, [14]- [16], and the generalized Lagrange metrics on M, given by

$$h_{ij} = \frac{1}{12F^4} \frac{\partial^2 F^4}{\partial y^i \partial y^j}, \quad \tilde{h}_{ij} = \frac{1}{12F^6} \frac{\partial^2 F^4}{\partial y^i \partial y^j}.$$
(9)

(*h* defined above is the same as the one in [5], with the only difference that here we have divided by F^4 or F^6 instead of F^2 , in order that the obtained tensors be homogeneous of degree -2, respectively, -4).

In the following, we shall use two particular kinds of metrics on $\widetilde{T^2M}$, namely:

1.
$$g_{ij} = g_{ij} = g_{ij}(x), \quad g_{ij} = \tilde{h}_{ij}(y^{(1)}),$$

2. $g_{ij} = g_{ij}(x), \quad g_{ij} = g_{ij} = h_{ij}(y^{(1)}),$

 $g_{ij}(x)$ being a Riemannian metric on M, and h_{ij} , \tilde{h}_{ij} as above.

These two examples have an important property, namely, they are compatible to the almost contact structures \mathbb{F} introduced in [1].

An N-linear connection D is called *metrical* if $D_X G = 0$, $\forall X \in \mathcal{X}(T^2 M)$. The local expression of this equality is given in [1].

4 The Ricci tensor Ric(D)

If we consider the Ricci tensor Ric(D), as the trace of the linear operator

$$V \mapsto R\left(V, X\right) Y, \,\forall V = V^{(0)i} \delta_i + V^{(1)i} \delta_{1i} + V^{(2)i} \delta_{2i} \in \mathcal{X}\left(T^2 M\right), \tag{10}$$

then, [3], the Ricci tensor Ric(D) has the following components:

$$\begin{aligned} \operatorname{Ric}\left(D\right)\left(\frac{\delta}{\delta x^{j}},\frac{\delta}{\delta x^{i}}\right) &= \operatorname{R}_{(00)}^{l}{}_{ijl}^{l} =: \operatorname{R}_{ij};\\ \operatorname{Ric}\left(D\right)\left(\frac{\delta}{\delta y^{(1)j}},\frac{\delta}{\delta x^{i}}\right) &= -\operatorname{P}_{(10)}^{l}{}_{ijj} =: -\operatorname{P}_{(10)}^{2}{}_{ij};\\ \operatorname{Ric}\left(D\right)\left(\frac{\delta}{\delta y^{(2)j}},\frac{\delta}{\delta x^{i}}\right) &= -\operatorname{P}_{(20)}^{l}{}_{ijl}^{l} =: -\operatorname{P}_{(20)}^{2}{}_{ij};\\ \operatorname{Ric}\left(D\right)\left(\frac{\delta}{\delta x^{j}},\frac{\delta}{\delta y^{(1)i}}\right) &= \operatorname{P}_{(11)}^{l}{}_{ijl}^{l} =: \operatorname{P}_{(11)}^{1}{}_{ij};\\ \operatorname{Ric}\left(D\right)\left(\frac{\delta}{\delta y^{(2)j}},\frac{\delta}{\delta y^{(1)i}}\right) &= \operatorname{S}_{(11)}^{l}{}_{ijl}^{l} =: \operatorname{S}_{(1)}^{i}{}_{ij};\\ \operatorname{Ric}\left(D\right)\left(\frac{\delta}{\delta y^{(2)j}},\frac{\delta}{\delta y^{(1)i}}\right) &= -\operatorname{Q}_{(21)}^{l}{}_{ijl}^{l} =: -\operatorname{Q}_{(21)}^{2}{}_{ij};\\ \operatorname{Ric}\left(D\right)\left(\frac{\delta}{\delta x^{j}},\frac{\delta}{\delta y^{(2)i}}\right) &= \operatorname{P}_{(22)}^{l}{}_{ijl}^{l} =: \operatorname{P}_{(22)}^{1}{}_{ij};\end{aligned}$$

$$\begin{aligned} \operatorname{Ric}\left(D\right)\left(\frac{\delta}{\delta y^{(1)j}},\frac{\delta}{\delta y^{(2)i}}\right) &= \begin{array}{c} Q_{i \ j l}^{l} \ =: \begin{array}{c} 1\\ Q_{i j l} \end{array};\\ \operatorname{Ric}\left(D\right)\left(\frac{\delta}{\delta y^{(2)j}},\frac{\delta}{\delta y^{(2)i}}\right) &= \begin{array}{c} S_{i \ 22}^{l} \ =: \begin{array}{c} S_{i j l} \end{array};\end{aligned}$$

5 Canonical structures

Let (M, g) be a Riemannian manifold and T^2M , its second order tangent bundle. The canonical nonlinear connection N is defined (cf. with R. Miron and Gh. Atanasiu, [13]) by its dual coefficients

$$M_{(1)j}^{i} = \gamma_{jk}^{i} y^{(1)k}, \quad M_{(2)j}^{i} = \frac{1}{2} \left\{ \mathbb{C} \left(\gamma_{jk}^{i} y^{(1)k} \right) + M_{(1)k}^{i} M_{(1)j}^{k} \right\},$$
(11)

 $\gamma_{jk}^i = \gamma_{jk}^i(x)$ being the Christoffel symbols of g and $\mathbb{C} = y^{(1)i} \frac{\partial}{\partial x^i} + 2y^{(2)i} \frac{\partial}{\partial y^{(1)i}}$. Let

$$\underset{(1)}{N^{i}}{}^{j} = \underset{(1)}{M^{i}}{}^{j}, \ \underset{(2)}{N^{i}}{}^{j} = \underset{(2)}{M^{i}}{}^{j} + \underset{(1)}{M^{i}}{}^{k} \underset{(1)}{M^{k}}{}^{k}$$

be its (direct) coefficients. Then, the coefficients of the Lie brackets, [1],

$$\delta_{0j}, \delta_{0k}] = \underset{(01)}{R} \underset{(12)}{^{i}} \delta_{1i} + \underset{(02)}{R} \underset{(21)}{^{i}} \delta_{2i}, \quad [\delta_{0j}, \delta_{1k}] = \underset{(11)}{B} \underset{(11)}{^{i}} \delta_{1i} + \underset{(12)}{B} \underset{(21)}{^{i}} \delta_{2i}$$

$$[\delta_{0j}, \delta_{2k}] = \underset{(21)}{B} \underset{(21)}{^{i}} \delta_{1i} + \underset{(22)}{B} \underset{(21)}{^{i}} \delta_{2i}, \quad [\delta_{1j}, \delta_{1k}] = \underset{(12)}{R} \underset{(12)}{^{i}} \delta_{2i}$$

$$[\delta_{1j}, \delta_{2k}] = \underset{(21)}{B} \underset{(21)}{^{i}} \delta_{2i}, \quad [\delta_{2j}, \delta_{2k}] = 0$$
(12)

have the property that

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$$B^{i}_{(11)jk} = B^{i}_{(22)jk} = \gamma^{i}_{jk}, B^{i}_{(21)jk} = R^{i}_{(12)jk} = R^{i}_{(22)jk} = 0.$$
 (13)

In this paper, we shall use the metrical N-linear connection introduced by the first author, [1], given by the coefficients:

$$\begin{array}{ll}
 L_{(00)}^{i}{}_{jk} &= \frac{1}{2} g^{il} (\delta_{k} g_{jl} + \delta_{j} g_{lk} - \delta_{l} g_{jk}) \\
 L_{(\beta0)}^{i}{}_{jk} &= B_{(\beta\beta)}^{i}{}_{kj}^{i} + \frac{1}{2} g^{il} (\delta_{k} g_{jl} - B_{(\beta\beta)}^{m}{}_{kj}^{g} g_{ml} - B_{(\beta\beta)}^{m}{}_{kl}^{g} g_{jm}) \\
 C_{(\beta1)}^{i}{}_{jk} &= \frac{1}{2} g^{il} \delta_{1k} g_{jl}, \ (\delta = 0, 2), \\
 C_{(\epsilon2)}^{i}{}_{jk} &= \frac{1}{2} g^{il} \partial_{2k} g_{jl}, \ (\epsilon = 0, 1), \\
 C_{(\beta\beta)}^{i}{}_{jk} &= \frac{1}{2} g^{il} (\delta_{\beta k} g_{jl} + \delta_{\beta j} g_{lk} - \delta_{\beta l} g_{jk}), \ \delta_{2i} = \partial_{2i},
\end{array}$$
(14)

where $\beta = 1, 2$.

Then, we have to remark that, taking into account the relations (13), two of the coefficients of the torsion tensor vanish, namely

$$P_{(21)}^{i}{}^{j}_{jk} = S_{(12)}^{i}{}^{j}_{jk} = 0, (15)$$

where $P_{(21)}^{i}{}_{jk}\delta_{1i} = v_1 T(\delta_{2k}, \delta_j), \ S_{(12)}^{i}{}_{jk}\dot{\partial}_{2i} = v_2 T(\delta_{1k}, \delta_{1j}).$

6 The case of the g - h - h-metric

Let the metric structure of T^2M be given by

$$G = g_{ij}(x)dx^i \otimes dx^j + h_{ij}(y^{(1)})\,\delta y^{(1)i} \otimes \delta y^{(1)j} + h_{ij}(y^{(1)})\delta y^{(2)i} \otimes \delta y^{(2)j}$$

where g is a Riemannian metric on M and h is as in (9). Then, G is h-Riemannian and v_1 -, v_2 -locally Minkovski. In this case, the detailed expressions of the coefficients $D\Gamma(N)$ of the canonical N-linear connection and of its curvatures and torsions are given in [1].

By applying the results in the cited paper and the relation (15), we obtain by a direct computation the following result:

Proposition 1. The only nonvanishing components of the Ricci tensor Ric(D) of the canonical -linear connection are

$$Ric(D)\left(\frac{\delta}{\delta x^{j}},\frac{\delta}{\delta x^{i}}\right) = \underset{(00)^{i}jk}{Ric(D)} = \underset{(11)^{i}jk}{Ric(D)} = \underset{(11)^{i}jk}{P} = \underset{(11)^{i}jk}{P} = \underset{(11)^{i}jk}{P} = \underset{(11)^{i}jk}{P}$$
$$Ric(D)\left(\frac{\delta}{\delta y^{(1)j}},\frac{\delta}{\delta y^{(1)i}}\right) = \underset{(11)^{i}jk}{S} = \underset{(1)^{i}jk}{S} = \underset{(1)^{i}jk}{S}$$

where $r_{ij} = r_{i \ jk}^k$ denotes the Ricci tensor of the Levi-Civita connection attached to g.

By applying the results in [3], we can state:

Proposition 2. The Einstein equations associated to the metrical N-linear connection D are

$$R_{ij} - \frac{1}{2}(r + S_{(1)})g_{ij} = \kappa \mathcal{T}_{(00)_{ij}};$$

$$P_{(11)}^{1} = \kappa \mathcal{T}_{(10)}i_{j};$$

$$S_{(1)}^{1} - \frac{1}{2}(r + S_{(1)})h_{ij} = \kappa \mathcal{T}_{(11)}i_{j}, \alpha = 1, 2;$$

$$\mathcal{T}_{(20)}^{1} = \mathcal{T}_{(01)}^{1}i_{j} = \mathcal{T}_{(21)}i_{j} = \mathcal{T}_{(02)}i_{j} = \mathcal{T}_{(12)}i_{j} = \mathcal{T}_{(22)}i_{j} = 0.$$

7 The case of the $g - g - \tilde{h}$ -metric

Proposition 3. Now, let the metric structure of $\widetilde{T^2M}$ be given by

$$G = g_{ij}(x)dx^i \otimes dx^j + g_{ij}(x)\,\delta y^{(1)i} \otimes \delta y^{(1)j} + \widetilde{h}_{ij}(y^{(1)})\delta y^{(2)i} \otimes \delta y^{(2)j},$$

where g is a Riemannian metric on M and \tilde{h} is as in (9). Then, G is h-, v_1 - Riemannian and v_2 -locally Minkovski.

In order to determine the components of the Ricci tensor, we first have to compute the coefficients of the canonical N-linear connection in our case. We have:

Using the expressions above, we obtain

Proposition 4. All the components of the Ricci tensor of the N-linear connection D vanish, except

$$Ric(D)\left(\frac{\delta}{\delta x^{j}},\frac{\delta}{\delta x^{i}}\right) = \underset{(00)}{R}^{l}{}_{jl} =: r_{ij},$$

where r_{ij} denotes the Ricci tensor of the Levi-Civita connection of the metric g on M.

As a consequence, the Einstein equations can be written in this case as:

$$r_{ij} - \frac{1}{2}rg_{ij} = \kappa \mathop{\mathcal{T}}_{(00)}{}_{ij},$$

the other components of the energy-momentum tensor being identically 0. The equations above are exactly the Einstein equations of the Levi-Civita connection ∇ of g = g(x). Obviously, the energy conservation law is satisfied.

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